

Ch 6 习题: 本节假设 L -段有界域 $U \subset \mathbb{R}^n$ 为有界开集. $\partial U \in C^\infty$.

[6.1] 考虑带位势 C 的 Laplace 方程 $-\Delta u + Cu = 0 \dots (*)$

和散度形式的方程 $-\operatorname{div}(a \nabla v) = 0 \quad a > 0$.

(1) 证明: 若 u 为 $(*)$ 的解, $w > 0$ 也是 $(*)$ 的解, 则 $v = \frac{u}{w}$ 是 $(**)$ 的解 ($a = w^2$)

(2) 反之, 若 v 是 $(**)$ 的解, 则 $u = va^{\frac{1}{2}}$ 是 $(*)$ 的解, (对某个位势 C).

证明: (1). $-\Delta u + Cu = 0, \quad -\Delta w + Cw = 0$

$$v = \frac{u}{w}$$

$$\Rightarrow \partial_i v = \frac{\partial_i u \cdot w - u \cdot \partial_i w}{w^2} \Rightarrow a \partial_i v = \frac{\partial_i u \cdot w - u \cdot \partial_i w}{w^2}$$

$$-\operatorname{div}(a \nabla v) = -\sum_{i=1}^n \partial_i (a \partial_i v)$$

$$= -\sum_{i=1}^n \partial_i a \partial_i v - \sum_{i=1}^n a \partial_i \partial_i v$$

$$= -\sum_{i=1}^n \partial_i a \frac{\partial_i u \cdot w - u \cdot \partial_i w}{w^2} - \sum_{i=1}^n a \frac{\partial_i (\partial_i u \cdot w - u \cdot \partial_i w) w^2 - 2w \partial_i w (\partial_i u \cdot w - u \cdot \partial_i w)}{w^4}$$

$$a = w^2 \Rightarrow a \partial_i v = \partial_i u \cdot w - \partial_i w \cdot u$$

$$\operatorname{div}(a \nabla v) = \sum_{i=1}^n \partial_i (\partial_i u \cdot w - \partial_i w \cdot u)$$

$$= \sum_{i=1}^n (\partial_i \partial_i u \cdot w + \partial_i u \partial_i w - \partial_i w \partial_i u - \partial_i \partial_i w \cdot u)$$

$$= \Delta u \cdot w - \Delta w \cdot u$$

$$= Cw - Cw = 0.$$

(2) 若 $-\operatorname{div}(a \nabla v) = 0$.

$$\text{则 } \sum_{i=1}^n \partial_i (a \partial_i v) = 0 \Rightarrow \sum_{i=1}^n \partial_i a \partial_i v + a \sum_{i=1}^n \partial_i \partial_i v = 0$$

$$-\Delta u + Cu = -\sum_{i=1}^n \partial_i (\partial_i (va^{\frac{1}{2}})) + Cva^{\frac{1}{2}}$$

$$v = a^{\frac{1}{2}} v$$

$$= Cva^{\frac{1}{2}} - \sum_{i=1}^n \partial_i (\partial_i v \cdot a^{\frac{1}{2}} + \frac{1}{2} a^{\frac{1}{2}} \partial_i a \cdot v)$$

$$= Cva^{\frac{1}{2}} - \sum_{i=1}^n (\partial_i \partial_i v \cdot a^{\frac{1}{2}} + \partial_i v \cdot \frac{1}{2} a^{\frac{1}{2}} \partial_i a + \frac{1}{2} a^{\frac{1}{2}} \partial_i a \partial_i v + \frac{1}{4} (\partial_i \partial_i a^{\frac{1}{2}}) v)$$

$$= Cva^{\frac{1}{2}} - a^{-\frac{1}{2}} \cdot (\sum_{i=1}^n \partial_i a \partial_i v + a \sum_{i=1}^n \partial_i \partial_i v) - \frac{1}{4} \sum_{i=1}^n (\partial_i \partial_i a^{\frac{1}{2}}) v$$

$$= \frac{1}{4} v (Ca^{\frac{1}{2}} - \Delta \sqrt{a})$$

$$\text{取 } c = \frac{\Delta \sqrt{a}}{\sqrt{a}} \quad \text{即证} \quad \square$$

[6.2]. 设 $Lu = -\sum_{i,j=1}^n a^{ij} \partial_{ij} u + cu$.

证明: 存在常数 $\mu > 0$, s.t. $\mathbb{R}^n \times \mathbb{R} \ni (x, u) \rightarrow -\mu(x, u)$ 的条件下, $B[\cdot, \cdot]$ 满足 (or Milgram 定理条件)

证明: $B[u, v] = \sum_{i,j=1}^n \int_{\Omega} a^{ij} \partial_{ij} u \partial_{ij} v + cuv \quad \forall u, v \in H_0^1(\Omega)$

$$\textcircled{1} |B[u, v]| \leq \|a^{ij}\|_{L^\infty} \sum_{i,j=1}^n \int_{\Omega} |\partial_{ij} u| |\partial_{ij} v| + \|c\|_{L^\infty} \int_{\Omega} |u| |v| dx$$

$$\stackrel{\text{Hölder}}{\leq} C \left(\|Du\|_{L^2} \|Dv\|_{L^2} + \|u\|_{L^2} \|v\|_{L^2} \right)$$

$$\leq C (\|u\|_{H_0^1} \|v\|_{H_0^1})$$

$$\textcircled{2} |B[u, u]| = \int_{\Omega} \sum_{i,j=1}^n a^{ij} \partial_{ij} u \partial_{ij} u + cu^2$$

$$\stackrel{L\text{-致密性}}{\geq} \theta \|Du\|_{L^2}^2 + \int_{\Omega} cu^2$$

~~$u \in H_0^1$, 由 Poincaré 不等式 $\|u\|_{L^2} \leq C' \|Du\|_{L^2}$ (for some $C' > 0$).~~

~~$$= \theta \|Du\|_{L^2}^2 + (c + \mu) \int_{\Omega} u^2 - (\mu + \epsilon) \int_{\Omega} u^2$$~~

~~Poincaré 不等式: $\forall u \in H_0^1(\Omega)$, 故 $\exists C' > 0, \|u\|_{L^2} \leq C' \|Du\|_{L^2}$~~

~~$$\rightarrow (\theta - c^2 \mu^{-2} \epsilon) \|Du\|_{L^2}^2 + (c + \mu) \int_{\Omega} u^2$$~~

~~$$\text{取 } \mu + \epsilon \geq \theta - c^2 \mu^{-2} \epsilon \geq C_0, \text{ 即 } \mu \leq \frac{\theta - (1 + \frac{1}{c^2}) \epsilon_0}{\epsilon_0}$$~~

~~$$\rightarrow \geq \theta \|Du\|_{L^2}^2 - C(C')^2 \|Du\|_{L^2}^2$$~~

$$= \theta \|Du\|_{L^2}^2 + \int_{\Omega} (c + \mu + \epsilon) u^2 - (\mu + \epsilon) \int_{\Omega} u^2$$

Poincaré: $\forall u \in H_0^1(\Omega)$, then $\exists C' > 0, \|u\|_{L^2} \leq C' \|Du\|_{L^2}$.

$$\geq \theta \|Du\|_{L^2}^2 + \int_{\Omega} (c + \mu + \epsilon) u^2 - (\mu + \epsilon) (C')^2 \|Du\|_{L^2}^2$$

$$\text{取 } (\mu + \epsilon)(C')^2 = \frac{\theta}{2} \quad (\epsilon \text{ 足够小}), \text{ 于是}$$

$$\text{上式} \geq \frac{\theta}{2} \|Du\|_{L^2}^2 = \frac{\theta}{4} \|Du\|_{L^2}^2 + \frac{\theta}{4} \|Du\|_{L^2}^2$$

$$\stackrel{\text{Poincaré}}{\geq} \frac{\theta}{4} \|Du\|_{L^2}^2 + \frac{\theta}{4} \frac{1}{C'} \|u\|_{L^2}^2$$

$$\geq \min \left\{ \frac{\theta}{4}, \frac{\theta}{4} \frac{1}{C'} \right\} \|u\|_{H_0^1}^2$$

□

(6.3) $u \in H_0^2(U)$ 是如下边值问题 $\begin{cases} \Delta^2 u = f & \text{in } U \\ u = \frac{\partial u}{\partial \nu} = 0 & \text{on } \partial U \end{cases}$ 的弱解, 若 $\int_U \Delta u \Delta v \, dx = \int_U f v \, dx \quad \forall v \in H_0^2(U)$

今给定 $f \in L^2(U)$, 证明该方程存在唯一-弱解

证明: 令 $B[u, v] = \int_U \Delta u \Delta v \, dx$

(1) $|B[u, v]| = \left| \int_U \Delta u \Delta v \, dx \right|$
 $\stackrel{\text{Holder}}{\leq} C \|\Delta u\|_{L^2} \|\Delta v\|_{L^2}$

$\leq C' \|u\|_{H_0^2} \|v\|_{H_0^2}$
 $u, v \in H_0^2(U)$
 (由 Poincaré 知 $\|u\|_2, \|\Delta u\|_2$ 由 $\|\Delta^2 u\|_2$ 控制)

(2) $B[u, u] \stackrel{\text{若}}{=} \int_U \Delta u \Delta u \, dx$
 $= \sum_{j,k=1}^n \int \partial_j^2 u \partial_k^2 u \, dx$

$\stackrel{\text{分部积分}}{=} - \sum_{j,k=1}^n \int \partial_j u \cdot \partial_j \partial_k^2 u \, dx$

$\stackrel{\text{再分部积分}}{=} \sum_{j,k=1}^n \int (\partial_j \partial_k u)^2 \, dx = \|\Delta u\|_{L^2}^2$

$\stackrel{\text{Poincaré}}{\geq} C (\|\Delta^2 u\|_{L^2} + \|u\|_{L^2} + \|\Delta u\|_{L^2})^2$
 $= C \|u\|_{H_0^2(U)}^2$

~~对 $\forall u \in H_0^2(U)$ 可以找到 $u_n \in C_c^\infty(U)$~~

~~s.t. $\|u_n - u\|_{H_0^2(U)} \rightarrow 0$ as $n \rightarrow \infty$~~

对 $\forall u \in H_0^2(U)$, $\exists \{u_n\} \subset C_c^\infty(U)$ s.t. $\|u_n - u\|_{H_0^2(U)} \rightarrow 0$

故 $\|u_n\|_{H_0^2} \rightarrow \|u\|_{H_0^2}$

$|\|\Delta u_n\|_{L^2}^2 - \|\Delta u\|_{L^2}^2| \leq \|\Delta(u_n - u)\|_{L^2}^2 \leq C \|\Delta^2(u_n - u)\|_{L^2}^2 \rightarrow 0$

$\therefore B[u, u] \geq C \|u\|_{H_0^2(U)}^2$ 对 $u \in H_0^2(U)$ 成立

即由 (1)(2), 据 Lax-Milgram 定理知 $\exists! u \in H_0^2(U)$

s.t. $\forall v \in H_0^2(U)$, $B[u, v] = (f, v)_{L^2}$ given $f \in L^2$.

□

4. 设 U 是连通的, $u \in H^1(U)$ 是 $\begin{cases} -\Delta u = f & \text{in } U \\ \frac{\partial u}{\partial n} = 0 & \text{on } \partial U \end{cases}$ 的弱解. 若 $\int_U Du \cdot Du \, dx = \int_U f u \, dx$

设 $f \in L^2(U)$. 证明该方程弱解存在 $\Leftrightarrow \int_U f \, dx = 0$ $\forall u \in H^1(U)$

证明: \Rightarrow : 反证: 若 $\int_U f \, dx \neq 0$.

则设 u 为该方程的弱解. (因 U 有界)

$$\int_U Du \cdot Du = \int_U Du \cdot D(v - \langle v \rangle) \stackrel{\substack{v - \langle v \rangle \in H^1(U) \\ \text{由弱解定义}}}{=} \int_U f(v - \langle v \rangle)$$

$$\langle v \rangle = \frac{1}{|U|} \int_U v \, dx \in \mathbb{R}.$$

$$= \int_U f v - \langle v \rangle \int_U f \neq \int_U f v$$

这与 $\forall v \in H^1(U), \int_U Du \cdot Du = \int_U f v$ 矛盾!

只要取 $\langle v \rangle \neq 0$ 即可

于是 $\int_U f \, dx = 0$

\Leftarrow : 若 $\int_U f \, dx = 0$.

设 $H_0^1(U) = \{v \in H^1(U) \mid \int_U v \, dx = 0\}$ Hilbert 空间

内积 $(u, v)_{H_0^1} = \int_U Du \cdot Dv \, dx$

$F(v) = \int_U f v \, dx : H_0^1(U) \rightarrow \mathbb{R}$

$v \mapsto \int_U f v \, dx$ 为 $H_0^1(U)$ 上的连续线性泛函.

由 Riesz 表示定理, 存在唯一 $u \in H_0^1(U)$ s.t.

$F(v) = (u, v)_{H_0^1} \Rightarrow \int_U Du \cdot Dv = \int_U f v$ 从而得证

check: $H_0^1(U)$ 是 Hilbert 空间;

且用 check ①: $(u, u)_{H_0^1} = 0 \iff u = 0$.

② $\|\cdot\|_{H_0^1}$ 与 $\|\cdot\|_{H^1}$ 等价. 这由 Poincaré 不等式易得.

这是因为 $(u, u)_{H_0^1} = \|Du\|_{L^2}^2 = 0 \Rightarrow Du = 0$ a.e. 又 $u \in H_0^1(U)$ 故 $\langle u \rangle = 0 \Rightarrow u = 0$ a.e.



[5]. $\begin{cases} -\Delta u = f & \text{in } U \\ u + \frac{\partial u}{\partial n} = 0 & \text{on } \partial U \end{cases}$ 问: 如何定义 H^1 弱解?
 给定 $f \in L^2(U)$ 时, 如何证明解的存在唯一性?

Proof: 令 $B[u, v] = \int_U Du \cdot Dv \, dx + \int_{\partial U} \text{Tr} u \cdot \text{Tr} v \, d\mathcal{H}^{n-1}$. $u, v \in H^1(U)$.

这么定义是因为, 假设 $u, v \in C^\infty(U)$, 则

$$\begin{aligned} \int_U -\Delta u \cdot v &\stackrel{\text{Gauss 公式}}{=} \int_U Du \cdot Dv - \int_{\partial U} v Dn \cdot \vec{\nu} \, d\mathcal{H}^{n-1} \\ &= \int_U Du \cdot Dv - \int_{\partial U} v \frac{\partial u}{\partial \nu} \, d\mathcal{H}^{n-1} \\ &= \int_U Du \cdot Dv + \int_{\partial U} u \cdot v \, d\mathcal{H}^{n-1}. \end{aligned}$$

(1) $|B[u, v]| \leq C \|Du\|_{H^1} \|v\|_{H^1}$ 是显然的 (由上述定理易得)

(2) 反之: $B[u, u] = \int_U Du \cdot Du \, dx + \int_{\partial U} (\text{Tr} u)^2 \, d\mathcal{H}^{n-1}$.
 $\exists \beta > 0$
 $\geq \beta \|u\|_{H^1(U)}^2 \quad \forall u \in H^1(U)$

反之: 若不成立, 则 $\forall n \in \mathbb{Z}_+$ $\exists u_n$. $\|u_n\|_{H^1(U)}^2 = 1$

$$n B[u_n, u_n] < \frac{1}{n} \|u_n\|_{H^1(U)}^2 = 1.$$

$$\text{由 } B[u_n, u_n] < \frac{1}{n}$$

由于 $\{u_n\} \subset H^1(U)$ - 没有界. 由 Banach-Alaoglu 定理,

\exists 子列 $u_{n_k} \rightharpoonup$ for some $u \in H^1(U)$ in $H^1(U)$.

由 $H^1(U) \hookrightarrow L^2(U)$ (Rellich-Kondrachev) 故

$$u_{n_k} \rightarrow u \text{ in } L^2(U).$$

$$\text{但 } \|Du_{n_k}\|_{L^2}^2 \leq B[u_{n_k}, u_{n_k}] \leq \frac{1}{n_k} \rightarrow 0 \text{ as } k \rightarrow \infty$$

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5月6日 习题课.

[6.6] Ω 连通, $\partial\Omega$ 由两个不交闭集 Γ_1, Γ_2 构成. 请定义如下方程的弱解.

$$\begin{cases} -\Delta u = f & \text{in } \Omega \\ u = 0 & \text{on } \Gamma_1 \\ \frac{\partial u}{\partial \nu} = 0 & \text{on } \Gamma_2 \end{cases} \quad \text{并讨论存在、唯一性}$$

证明: $H_0^1(\Omega) := H^1(\Omega) \cap \{u \mid \text{Tr } u|_{\Gamma_1} = 0\}$

在 $H_0^1(\Omega)$ 上定义 $B[u, v] = \int \nabla u \cdot \nabla v$.

则 $H_0^1(\Omega)$ 是 $H^1(\Omega)$ 的闭子空间

* 称 $u \in H_0^1(\Omega)$ 是原方程的弱解, 若 $\forall v \in H_0^1(\Omega), B[u, v] = \int_{\Omega} f v \, dx$

为什么? 形式上: 对 C^∞ 函数

$$\int_{\Omega} -\Delta u \cdot v = \int_{\Omega} f v$$

$$\int_{\Omega} -\nabla u \cdot \nabla v \, dx + \int_{\partial\Omega} -\nabla u \cdot \nu$$

$$= \int_{\Omega} \nabla u \cdot \nabla v - \int_{\partial\Omega} \frac{\partial u}{\partial \nu} \cdot \nu \quad \begin{cases} \Gamma_1 & u = 0 \leftarrow \text{这步与应用 } u \in C_c^\infty \text{ 有关} \\ \Gamma_2 & \frac{\partial u}{\partial \nu} = 0 \leftarrow \text{条件} \end{cases}$$

剩下与 5 的证明类似. (Lax-Milgram). □

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Fredholm = 择一与椭圆算子特征值问题:

• 择一: Lax-Milgram 定理只回答了 $L + \mu \cdot I$ 在 μ 较大时, 对应椭圆方程 $\exists!$ 的 H_0^1 弱解, 但一般情况如何, 不得而知.

事实上, 一般情况会出现 "择一"; 这导致代类的

Given $f \in L^2(U)$

$$\begin{cases} Lu = f & \text{in } U \\ u = 0 & \text{on } \partial U \end{cases} \quad \exists! \text{ 弱解.}$$

Case 1: $\exists!$ 弱解 $\Leftrightarrow \forall v \in N^* \langle f, v \rangle = 0$

Case 2: 存在非零弱解

$$\begin{cases} Lu = 0 & \text{in } U \\ u = 0 & \text{on } \partial U \end{cases} \quad \text{解空间 } N \subset H_0^1(U) \text{ 满足} \\ \dim N = \dim N^* < +\infty$$

$$N^* = \left\{ \begin{cases} L^* v = 0 & \text{in } U \\ v = 0 & \text{on } \partial U \end{cases} \right. \quad \text{的解空间}$$

定义: $B_V[u, v] = B[u, v] + \nu(u, v)$

$$L_V u = Lu + \nu u$$

$$\forall g \in L^2, \exists! u \in H_0^1(U), B_V[u, v] = (g, v), \forall v \in H_0^1(U)$$

$$u = L_V^{-1} g = L_V^{-1}(\nu u + f) \xrightarrow{k \Rightarrow L_V^{-1}} u(I - k)u = h = L_V^{-1} f$$

再用 Fredholm = 择一



在上述判之情形下, 特征值问题之解

$$\begin{cases} Lu = \lambda u + f & \text{in } U \\ u = 0 & \text{on } \partial U \end{cases}$$

• $\exists!$ 弱解 $\iff \lambda \notin \Sigma$ λ 不在 L 的谱里面即可

— $\lambda \in \Sigma$ 的例子: $U = (0, 2\pi) \times (0, 2\pi)$

$$\begin{cases} \Delta u + \frac{5}{4}u = ax_1 + bx_2 + c & \text{in } U \\ u = 0 & \text{on } \partial U \end{cases}$$

在此 U 上, 考虑

$$\begin{cases} \Delta u + \lambda u = 0 & \text{in } U \\ u = 0 & \text{on } \partial U \end{cases}$$

如何求解? \hookrightarrow 分离变量: $u = f(x_1)g(x_2)$

$$\text{代入} \Rightarrow f''g + fg'' + \lambda fg = 0, \quad \frac{f''}{f} + \frac{g''}{g} + \lambda = 0.$$

$$\text{设 } -s^2 = \frac{f''}{f}$$

$$-t^2 = \frac{g''}{g}$$

$$s^2 + t^2 = \lambda.$$

$$\begin{cases} f'' + s^2 f = 0 \\ f(0) = f(2\pi) = 0 \end{cases}$$

$$\begin{cases} g'' + t^2 g = 0 \\ g(0) = g(2\pi) = 0 \end{cases}$$

$$\Rightarrow f(x_1) = C_k \sin\left(\frac{k}{2}x_1\right).$$

$$g(x_2) = C_l \sin\left(\frac{l}{2}x_2\right)$$

$$\lambda = \frac{k^2 + l^2}{4}$$

\nwarrow 主特征值.

$$k, l = 1.$$

$$\lambda_1 = \frac{1}{2}$$

$$u_{01}(x_1, x_2) = \sin\frac{x_1}{2} \sin\frac{x_2}{2}$$

$$k=1$$

$$l=2$$

$$\lambda_2 = \frac{5}{4}$$

$$u_{12}(x_1, x_2) = \sin\frac{x_1}{2} \sin x_2$$

$$k=2$$

$$l=1$$

$$u_{21}(x_1, x_2) = \sin x_1 \sin\frac{x_2}{2}$$

\vdots

\vdots

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$$\begin{cases} \Delta u + \frac{5}{4} u = 0 & \text{in } U \\ u = 0 & \text{on } \partial U \end{cases}$$

解空间应由 u_{12}, u_{21} 张成

$$\therefore \text{原方程有解} \iff \begin{cases} (f, u_{12}) = 0 \\ (f, u_{21}) = 0 \end{cases} \quad f = ax_1 + bx_2 + c$$

$$\Rightarrow a = b = 0, \quad c \in \mathbb{R}$$

Moreover: $\forall v \in H_0^1(\Omega)$ 若 $\iint_{(0,2\pi) \times (0,2\pi)} v(x_1, x_2) \sin \frac{x_1}{2} \sin \frac{x_2}{2} dx_1 dx_2 = 0$

$$\text{则 } \|u\|_{L^2}^2 \leq \frac{4}{5} \|\nabla v\|_{L^2}^2$$

Pf: Consider:

$$\begin{cases} -\Delta w_{kl} = \lambda_k w_{kl} & \text{in } U \\ w_k = 0 & \text{on } \partial U \end{cases}$$

$$w_{kl} = \sin \frac{kx_1}{2} \sin \frac{l x_2}{2} \quad \{w_{kl}\} \text{ 为 } L^2 \text{ 的正交基}$$

$$v = \sum_{k,l=1}^{\infty} (v, w_{kl}) w_{kl}$$

$$\text{Fact: } - \int_U \Delta w_k \cdot w_p = \lambda_k \int_U w_k w_p dx = \lambda_k \|w_k\|^2$$

$$\text{从而 } \int_U |\nabla w_{kl}|^2 dx = \lambda_k \int_U w_{kl}^2 dx$$

$$\int_U \nabla w_k \cdot \nabla w_p dx = 0 \quad k \neq p$$

$$\begin{aligned} \therefore \|\nabla v\|_{L^2}^2 &= \int_U \left(\sum_{k,l} \alpha_{k,l} \nabla w_{kl} \right)^2 dx = \sum_{k,l} \alpha_{k,l}^2 \int_U |\nabla w_{kl}|^2 dx \\ &= \sum_{k,l} \alpha_{k,l}^2 \int_U \lambda_k w_{kl}^2 dx \geq \sum_{k,l} \frac{5}{4} \alpha_{k,l}^2 \int_U w_{kl}^2 dx \\ &= \frac{5}{4} \int_U (\sum_{k,l} \alpha_{k,l} w_{kl})^2 dx = \frac{5}{4} \int_U |v|^2 dx \end{aligned}$$

□

7. 设 $u \in H^1(\mathbb{R}^n)$ (具有紧支集) 是 $-\Delta u + c(u) = f$ in \mathbb{R}^n 的弱解.

其中 $f \in L^2(\mathbb{R}^n)$, $c: \mathbb{R} \rightarrow \mathbb{R}$ 是光滑函数, $c(0) = 0$, $c'(x) \geq 0$

$\Rightarrow c(u) \in L^2(\mathbb{R}^n)$. 证明: $\|D^2 u\|_2 \leq C \|f\|_2$

证明: u 是 $-\Delta u + c(u) = f$ 的 H^1 弱解.

则 $\forall v \in H^1(\mathbb{R}^n)$.

$$\int_U \nabla u \cdot \nabla v + c(u) \cdot v \, dx = \int_U f v \, dx$$

取 $v = -D_k^{-h} D_k^h u$, $0 < |h| \ll 1$. 则 $v \in H^1(\mathbb{R}^n)$ 且紧支.

代入:

$$\int_U D_u \cdot (-D_k^{-h} D_k^h u) \, dx + \int_U c(u) D_k^{-h} D_k^h u \, dx = - \int_U f D_k^{-h} D_k^h u \, dx$$

• D 与 D_k^h 可交换.

• 差商的形式分部积分

$$\Rightarrow \int_U |D_k^h D_k^h u|^2 + \underbrace{D_k^h c(u) \cdot D_k^h u}_{\text{希望估计的是这项}} \, dx = - \int_U f \cdot D_k^{-h} D_k^h u \, dx$$

\rightarrow 因此要先设法估计非线性项 $c(u)$ 带来的贡献.

$$\phi \frac{D_k^h c(u)(x) D_k^h u(x)}{h} \phi = \left(\frac{c(u(x+h e_k)) - c(u(x))}{h} \right) \cdot D_k^h u(x) \phi$$

中值定理 $\exists \xi \in \mathbb{R}$

$$= \phi c'(\xi) \cdot \frac{|u(x+h e_k) - u(x)|}{h} |D_k^h u(x)|$$

$$= \phi c'(\xi) \cdot |D_k^h u(x)|^2 \geq 0$$

$$\Rightarrow \int_U |D_k^h D_k^h u|^2 \leq - \int_U f D_k^{-h} D_k^h u \, dx \leq \left| \int_U f D_k^{-h} D_k^h u \right|^2$$

$$\leq C \|f\|_2^2 + \varepsilon \|D_k^{-h} D_k^h u\|_2^2$$

$$\leq C \|f\|_2^2 + C \varepsilon \|D_k^h D_k^h u\|_2^2$$

$\varepsilon < \frac{1}{2}$ 即可.
再用差商性质 \square

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椭圆方程正则性问题.

$$Lu = f \quad \text{in } U. \quad a, b, c \in C^{k,\alpha}(U), \quad f \in H^k(U)$$

$$\Rightarrow u \in H_{loc}^{k+2}(U). \quad (\text{符合常理}).$$

但未必所有方程都有此性质 (波方程, Schrodinger 方程 ^{形式上}).

椭圆正则性: 关注 "最高正则性"

Schrodinger: 低正则性问题

$$i\partial_t u + \Delta u = -|u|^2 u \quad \text{in } \mathbb{R}^3 \times \mathbb{R}$$

$$u|_{t=0} = u_0(x) \in H^s(\mathbb{R}^3), \quad s < 1.$$

之前: 局部解: $C_t^0 H^s \cap L_t^4 L_x^6$ (若初值 H^1).

global existence? $s > \frac{1}{3} \Rightarrow u \in C_t H_x^s(\mathbb{R} \times \mathbb{R}^3).$

(Bourgain: Scattering in Energy space & below for 3D NLS)

椭圆方程正则性证明方法: ① 用差商估计得到导数的估计.

② Bootstrap (高正则性):

$$\text{例如. } -\Delta u + \lambda u = f \quad \text{in } U. \quad \partial U \in C^\infty \quad f \in C^\infty$$

首先, 由 ~~Lax~~ 课本存在性理论, 某些 λ 使得 $\exists u \in H^1$ 作为弱解.

再用内正则性定理有 $\|u\|_{H^2} \lesssim \|\Delta u\|_2 + \|u\|_2.$

$$\text{此时. } \Delta u = \lambda u - f \in L^2 \Rightarrow \Delta u \in H^2$$

$$u \in H^2, f \in C^\infty \Rightarrow \text{Similarly } u \in H^4$$

$$\Rightarrow \dots \Rightarrow u \in H^{2k} \Rightarrow u \in C^\infty$$

对于解的可积性, 也有类似的 Bootstrap 方法

$$u \in H^1 = W^{1,2}, \quad \text{若 } d \geq 2.$$

$$\text{则由 Sobolev 嵌入定理 } W^{1,2} \hookrightarrow L^{2^*}, \quad 2^* = \frac{2d}{d-2} > 2$$

$$\Rightarrow \Delta u \in L^{2^*} \Rightarrow u \in W^{1,2^*}$$

\Rightarrow repeatedly and ~~and~~ derive ~~the~~ $u \in L^{\frac{d-\varepsilon}{2}}$ 不超过 d .

$$\frac{2d}{d-2} \downarrow$$

否则 Sobolev 嵌入

失效. \square

柯西估计在估计在处理一些其它方程时可能也会用到. 例如 Euler, Navier-Stokes 方程的湍流近稳定性, 往往需要用极高的正则性来换取长时间的稳定性. Bootstrap 仍适用. 但对柯西估计的具体估计也许会用到调和分析.

柯西方程的 ~~特征值问题~~ ~~一样~~:
极大值原理

弱极大值原理: L 想像成 " $-\Delta$ ".

强极大值原理: U 连通.

Hopf 引理: 结论很直观.

习题: 9. u 是 $L u = -\sum_{i,j=1}^n a^{ij} \partial_{ij} u = f$ in U 的 C^∞ 解. f bdd.
 $\begin{cases} u=0 & \text{on } \partial U. \end{cases}$

Fix $x^0 \in \partial U$. Define "barrier" at x_0 to be a C^2 function

$$w: L w \geq 1 \quad \text{in } U$$

证明: $\exists C > 0$

$$\left. \begin{array}{l} w(x^0) = 0 \\ w \geq 0 \end{array} \right\}$$

$$|\nabla u(x^0)| \leq C \left| \frac{\partial w}{\partial \nu}(x^0) \right|$$

$$w \geq 0 \quad \text{on } \partial U$$

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Proof: 对 w : 弱极大值原理, $\min_{\bar{U}} w = \min_{\partial U} w = w(x^0)$.

$$\text{令 } U_1 = u + w \|f\|_{\infty}$$

$$U_2 = u - w \|f\|_{\infty}$$

$$\text{则 } \Delta U_1 \geq 0$$

$$\Delta U_2 \leq 0$$

再由弱极大值原理: $\min_{\bar{U}} U_1 = \min_{\partial U} U_1 = \|f\|_{\infty} \min_{\partial U} w$
 $= \|f\|_{\infty} w(x^0) = U_1(x^0)$.

$$\max_{\bar{U}} U_2 = U_2(x^0)$$

由 Hopf 引理.

$$0 \geq \frac{\partial U_1}{\partial \nu}(x^0) = \frac{\partial u}{\partial \nu}(x^0) + \|f\|_{\infty} \frac{\partial w}{\partial \nu}(x^0)$$

$$\hookrightarrow \leq \frac{\partial U_2}{\partial \nu}(x^0) = \frac{\partial u}{\partial \nu}(x^0) - \|f\|_{\infty} \frac{\partial w}{\partial \nu}(x^0)$$

$$\text{又: } u=0 \text{ on } \partial U \quad \text{则 } \nabla u \parallel \bar{\nu}$$

$$\Rightarrow |\nabla u(x^0)| = \left| \frac{\partial u}{\partial \nu}(x^0) \right| \leq \|f\|_{\infty} \left| \frac{\partial w}{\partial \nu}(x^0) \right| \quad \square$$

10. (1) 连通 (2) 能量法 (3) 极大值原理

证明: $\begin{cases} -\Delta u = 0 & \text{in } U \\ \frac{\partial u}{\partial \nu} = 0 & \text{on } \partial U \end{cases}$ 的唯一光滑解为 $u = \text{const}$.

证明: (1) $I[w] = \frac{1}{2} \int_U |\nabla w|^2 dx$. $u = \text{const}$ 使 $I[w]$ 达到极小值 \checkmark

(2) 若 u 在 U 内达极大值, 那么由强极大值原理直接得 \checkmark

若 $\sup_{\bar{U}} u(x) = u(x^0) \quad x^0 \in \partial U$, 且 $\forall x \in U, u(x^0) > u(x)$

则 Hopf 引理表明: $\frac{\partial u}{\partial \nu}(x^0) > 0$. 矛盾! \square

$$12. Lu = -\sum_{i,j=1}^n a^{ij} \partial_{ij} u + \sum_{i=1}^n b^i \partial_i u + cu.$$

称 L 满足弱极大值原理, 是指 $\forall u \in C^2(U) \cap C(\bar{U})$.

今假设 $\exists v \in C^2(U) \cap C(\bar{U})$ 使 $Lv > 0$ in U , $v > 0$ on \bar{U} . 证明: L 满足弱极大值原理

Proof: 设 $u \in C^2(U) \cap C(\bar{U})$. $Lu \leq 0$ in U , $u \leq 0$ on ∂U .

~~$v \in C(U)$~~
 $\Rightarrow w = \frac{u}{v} \in C^2(U) \cap C(\bar{U})$.

希望构造一个有界算子 M . 使 $Mw \leq 0$ on $\{x \in \bar{U} \mid u > 0\} \subseteq U$.

如果能做到, 那么假设 $\{x \in \bar{U} \mid u > 0\}$ 非空.

由 M 满足弱极大值原理, 则有.

$$0 < \sup_{\{u > 0\}} w = \sup_{\{u > 0\}} w = \frac{0}{v} = 0. \quad \text{矛盾!}$$

下面构造 M . M 一致有界, 为达到此目的, 我们会保留一些项不变.

计算: $-a^{ij} \partial_{ij} w = -a^{ij} \partial_i \partial_j \left(\frac{u}{v} \right)$.

$$= -a^{ij} \partial_i \left(\frac{\partial_j u \cdot v - \partial_j v \cdot u}{v^2} \right) = \frac{-a^{ij}}{v^2}$$

$$= -a^{ij} \left(\partial_i \left(\frac{\partial_j u}{v} \right) - \partial_i \left(\frac{\partial_j v \cdot u}{v^2} \right) \right)$$

$$= -a^{ij} \left(\frac{\partial_i \partial_j u \cdot v - \partial_i u \partial_j v}{v^2} - \frac{-2v \partial_i v \partial_j u + v^2 \partial_i \partial_j v \cdot u + v^2 \partial_i v \partial_j u}{v^4} \right)$$

$$= \frac{-a^{ij} \partial_i \partial_j u \cdot v + a^{ij} \partial_i u \partial_j v}{v^2} + \frac{a^{ij} \partial_i v \partial_j u - a^{ij} \partial_i u \partial_j v}{v^2} - a^{ij} \frac{2}{v} \partial_i v \frac{\partial_j u \cdot v - \partial_j v \cdot u}{v^2}$$

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$$\sqrt{\frac{(Lu - \partial_i b^i u - cu)v + (-Lv + \partial_i b^i v + cv)u}{v^2}} + 0 + \frac{a^{ij} \frac{2}{v} \partial_j v \partial_i w}{v}$$

$$= \frac{Lu}{v} - \frac{uLv}{v^2} - b^i \partial_i w + a^{ij} \frac{2}{v} \partial_j v \partial_i w \quad (\text{上下指标表示求和}).$$

$$\text{令 } Mw = -a^{ij} \partial_{ij} w + \partial_i w (b^i - a^{ij} \frac{2}{v} \partial_j v)$$

$$= \frac{Lu}{v} - \frac{uLv}{v^2} \leq 0 \quad \text{on } \{x \in \bar{U} \mid u > 0\} \subset U$$

$$\begin{matrix} \uparrow \\ \text{为什么?} \end{matrix} \quad \begin{matrix} Lu \leq 0 \\ v > 0 \end{matrix} \quad \begin{matrix} u \leq 0 \\ Lv > 0 \end{matrix}$$

M-函数有性质。

□

一般的极大值原理。 $Lu = \partial_i (a^{ij} \partial_j u + b^i u) + c^i \partial_i u + du$

设 $|a^{ij}| \leq \lambda |\xi|^2$

$$\sum |a^{ij}(x)|^2 \leq \lambda^2$$

$$\frac{1}{\lambda^2} \sum (|b^i|^2 + |c^i|^2) + \frac{1}{\lambda} |d(x)| \leq v^2$$

Consider $Lu = g + \sum \partial_i f^i$ in U .

~~##~~

$$\left\{ \begin{array}{l} u = \varphi \\ \text{on } \partial U \end{array} \right.$$

$U \subset \mathbb{R}^d$ $f^i \in L^2(U)$. $g \in L^{q/2}(U)$. $q > d$.

u 为弱下解. $u \leq 0$ on ∂U . 则 $\sup_U \frac{u^+}{-u^-} \leq C (\|u^+\|_{L^2} + k)$

$$k = \frac{1}{\lambda} (\|f\|_{L^2} + \|g\|_{L^2})$$

证: De-Giorgi-Moser 迭代

证明思路: $\bar{u}^+ := u^+ + k$ (对弱下解).

$$H(z) = \begin{cases} z^\beta - k^\beta & k \leq z \leq N \\ \text{linear} & z \geq N \end{cases} \quad \underline{N > k}$$

只考虑 $d > 2$:

$$\text{先证: } \|H(\bar{u}^+)\|_{\frac{2d}{d-2}} \lesssim \left\| \frac{\cdot}{\cdot} \right\|_{\frac{2d}{d-2}}^{\frac{1}{2}} \| \bar{u}^+ \cdot H'(\bar{u}^+) \|_{\frac{2d}{d-2}}$$

$$q > 2^* = \frac{2d}{d-2}$$

下证 $N \rightarrow \infty$

$$\text{Fact: } \forall \beta \geq 1, \bar{u}^+ \in L^{\frac{2\beta d}{d-2}} \Rightarrow \bar{u}^+ \in L^{\frac{2\beta d}{d-2}}$$

$$q^* := \frac{2d}{d-2}, \quad \eta = \frac{d}{\frac{2d}{d-2}} \frac{2d/d-2}{2d/q-2} > 1$$

$$\|\bar{u}^+\|_{\beta \eta q^*} \lesssim \|\bar{u}^+\|_{\beta q^*}$$

$$\text{不断迭代: } \|\bar{u}^+\|_{\eta^N q^*} \lesssim \prod_{n=0}^{N-1} \|\bar{u}^+\|_{q^*}, \quad \bar{u}^+ \in \bigcap_{1 \leq p < \infty} L^p(\Omega)$$

$$N \rightarrow \infty \Rightarrow \|\bar{u}^+\|_{\infty} \lesssim \|\bar{u}^+\|_{q^*}$$

$$\Rightarrow \|\bar{u}^+\|_{\infty} \lesssim \|\bar{u}^+\|_{L^2}$$

□

具体参考

Gilbarg, Trudinger: Elliptic PDEs of 2nd order.
chapter 8.5 - 8.6.

local property. 更加精细.

能量法 (变分)

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思路: 方程的解 看作 能量泛函的极小值点.

难点: 给定方程, 如何找到能量泛函?

人为观察. \downarrow 要固定的

结论: 对 Laplace 方程:

结论:

$$\textcircled{1} \quad u \in C^2(\bar{U}) \text{ solves } \begin{cases} -\Delta u = f & \text{in } U \\ u = g & \text{on } \partial U \end{cases} \Leftrightarrow I(u) = \inf_{w \in \mathcal{A}} I[w]$$

$$\text{其中 } \mathcal{A} = \{w \in C^2(\bar{U}) \mid w = g \text{ on } \partial U\}$$

$$I[w] = \int_U \frac{1}{2} |\nabla w|^2 - f w \, dx.$$

$$\textcircled{2} \quad \begin{cases} -\Delta u = f \in C(\bar{U}) \\ \frac{\partial u}{\partial n} = g \in C(\partial U) \end{cases} \Leftrightarrow I(u) = \inf_{w \in \mathcal{A}} I[w]$$

$$\mathcal{A} = \left\{ w \mid \begin{array}{l} \frac{\partial w}{\partial n} = g \text{ on } \partial U \\ \int_U w \, dx = 0 \end{array} \right\}$$

$$I[w] = \int_U \frac{1}{2} |\nabla w|^2 - f w \, dx - \int_{\partial U} w g \, ds$$

$$\textcircled{3} \quad \begin{cases} -\Delta u = f \in C(\bar{U}) \\ \alpha(x)u + \frac{\partial u}{\partial n} = g \in C(\partial U) \end{cases}$$

$$u \text{ solves } \dots \Leftrightarrow I(u) = \inf_{w \in \mathcal{A}} I[w] \quad \text{其中 } \mathcal{A} = \{w \in C^2(\bar{U}) \cap C^1(\partial U)\}$$

$$I[w] = \int_U \frac{1}{2} |\nabla w|^2 - f w \, dx + \int_{\partial U} \left(\frac{1}{2} \alpha w \right)^2 - g w \, ds.$$

给出 (1) 的证明

⇒ 若 u 为方程解, $\forall w \in \mathcal{A}$

$$\begin{aligned}
 I[u] &= \frac{1}{2} \int_U |\nabla w|^2 - f w \, dx \\
 &= \frac{1}{2} \int_U \frac{|\nabla w|^2}{2} + \Delta u w \, dx \\
 &\stackrel{\text{分部}}{=} \frac{1}{2} \int_U |\nabla w|^2 + \sum_i \int_{\partial U} \partial_i (\partial_i u w) - \partial_i w \partial_i u \, dx \\
 &= \frac{1}{2} \int_U |\nabla w|^2 + \int_{\partial U} \nabla u \cdot w \cdot \vec{n} \, dS - \int_U \nabla u \cdot \nabla w \, dx \\
 &\geq \frac{1}{2} \int_U |\nabla w|^2 + \int_{\partial U} w \cdot \frac{\partial u}{\partial n} \, dS - \int_U \frac{|\nabla u|^2 + |\nabla w|^2}{2} \, dx \\
 &= -\frac{1}{2} \int_U |\nabla u|^2 \, dx + \int_{\partial U} u \cdot \frac{\partial u}{\partial n} \, dS \leftarrow \underline{u=g=u \text{ on } \partial U} \\
 &= -\frac{1}{2} \int_U |\nabla u|^2 \, dx + \int_U \partial_i (\partial_i u \cdot u) \, dx \\
 &= \frac{1}{2} \int_U |\nabla u|^2 \, dx - \int_U f u \, dx.
 \end{aligned}$$

⇐ : $\forall v \in C_c^\infty(U)$, $u+tv \in \mathcal{A} \quad \forall t$.

$i(t) = I[u+tv]$

$i(t)$ 在 $t=0$ 取极小. $\therefore i'(0) = 0$

⇒ $i'(t) = I'[u+tv] = \frac{1}{2} \int_U |\nabla u + t \nabla v|^2 \, dx - \int_U f u - \int_U t f v \, dx$

$i'(0) = 0$ 代入 $t=0$

$$\begin{aligned}
 0 &= \int_U \nabla u \cdot \nabla v - f v \, dx = \int_U (-\Delta u - f) v \, dx \\
 &\Rightarrow -\Delta u = f
 \end{aligned}$$

□

关于极大值原理, 有一套是构造神奇的辅助函数来应用

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例如: $u \in C^3(\Omega) \cap C^1(\bar{\Omega})$.

$$p. \begin{cases} \Delta u - 2u = f \\ u|_{\partial\Omega} = \varphi \end{cases} \quad \begin{array}{l} \Omega = B^1(0, 1) \subset \mathbb{R}^2 \\ f \in C(\bar{\Omega}), \varphi \in C(\partial\Omega). \end{array}$$

$$\text{则 } \sup_{\bar{\Omega}} |u| \leq \sup_{\partial\Omega} |\varphi| + \sup_{\bar{\Omega}} |f|$$

Proof: $F = \sup_{\bar{\Omega}} |f|, \Phi := \sup_{\partial\Omega} |\varphi|$

$$W := \Phi + \frac{1-x^2}{4} F \geq 0 \text{ in } \Omega.$$

$$\Delta W = -F \leq 0$$

$$\Rightarrow \Delta W - 2W = -F - 2W \leq -F \leq \pm f = \Delta(\pm u) - 2(\pm u)$$

\uparrow $W \geq 0$ \uparrow $F \geq 0$

$$W|_{\partial\Omega} = \Phi \geq \pm \varphi = \pm u|_{\partial\Omega}$$

$$\Rightarrow (\Delta - 2)(W \mp u) \leq 0 \text{ in } \Omega$$

$$W \mp u \geq 0 \text{ on } \partial\Omega$$

那么由极大值原理: $\pm u \leq W = \Phi + \frac{1-x^2}{4} F \leq \Phi + \frac{F}{4}$

□

2. 用于梯度估计

Global: $\Omega \subset \mathbb{R}^d$ 有界连通开. $u \in C^3(\Omega) \cap C^1(\bar{\Omega})$.

$$\Delta u = f(x, u) \in C^1(\Omega \times \mathbb{R}) \Rightarrow \sup_{\bar{\Omega}} |\nabla u| \leq C + \sup_{\partial\Omega} |\nabla u|$$

Proof: $\varphi = |\nabla u|^2 + \alpha u^2 + \beta \frac{|x|^2}{2d}, \alpha, \beta > 0$ 待定:

目标: $\Delta \varphi \geq 0$. 从而 $\sup_{\bar{\Omega}} |\nabla u| \leq \sup_{\partial\Omega} |\nabla u|$.

$$\text{取 } \sup_{\bar{\Omega}} |\nabla u|^2 \leq \sup_{\bar{\Omega}} \varphi \leq \sup_{\partial\Omega} \varphi = \sup_{\partial\Omega} |\nabla u|^2 + C$$

开方即可.

